

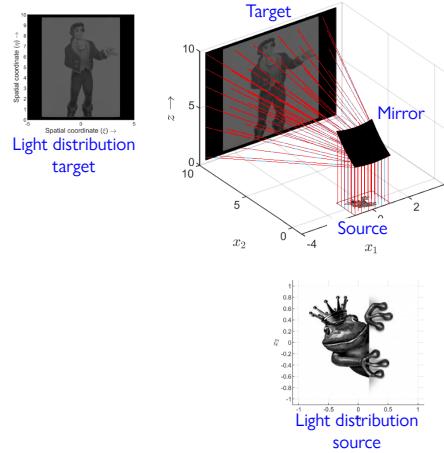
Computational illumination optics Wiskundige modellen voor optisch ontwerp

WND-conferentie, Noordwijkerhout, 15 december 2023

Martijn Anthonissen



Can we turn a frog into a prince?



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- ▶ "In the original Grimm version of the story, the frog's spell was broken when the princess threw the frog against the wall, at which he transformed back into a prince, while in modern versions the transformation is triggered by the princess kissing the frog."

https://en.wikipedia.org/wiki/The_Frog_Prince

- ▶ Can we do it with light?



Romijn, L. B. (2021)
Generated Jacobian Equations in Freeform Optical Design:
Mathematical Theory and Numerics
PhD thesis, Eindhoven University of Technology

Outline

Computational illumination optics at TU/e

Nonimaging freeform optics

Sixteen basic systems

1: Standard Monge-Ampère equation

2: Generalized Monge-Ampère equation

3: Generated Jacobian Equation

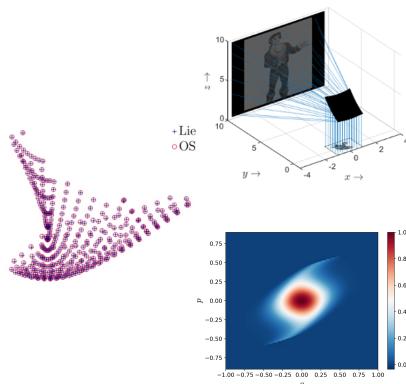
Iterative least-squares solver for 3D systems

Numerical results

Imaging optics

Improved direct methods

Conclusions and future work



Computational Illumination Optics Group at TU/e



Martijn Anthonissen



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Signify Research



Lisa Kusch



Koondi Mitra



Jan ten Thije
Boonkamp



Antonio Barion



Pieter Braam



Roel Hacking



Willem Jansen



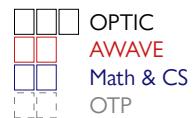
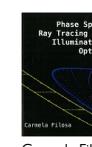
Vi Kronberg



Teun van Roosmalen



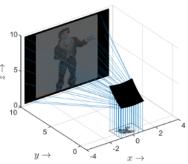
Sanjana Verma



Lines of research

Line A

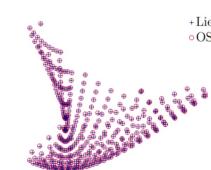
Nonimaging freeform optics



- ▶ Compute optical surfaces that convert given source into desired target distribution
- ▶ Freeform surfaces
- ▶ Fully nonlinear PDE of Monge-Ampère type

Line B

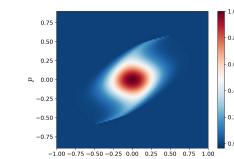
Imaging optics



- ▶ Make a very precise image of an object, minimizing aberrations
- ▶ Description with Lie transformations

Line C

Improved direct methods



- ▶ Ray tracing: iterative procedure to compute final design. Slow convergence
- ▶ Backward ray mapping
- ▶ Hamiltonian structure of equations and advanced numerical schemes for PDEs

Academic cooperation and industrial embedding

Design



Metrology



Manufacturing



Products



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Nonimaging freeform optics

Design of optical systems for illumination purposes

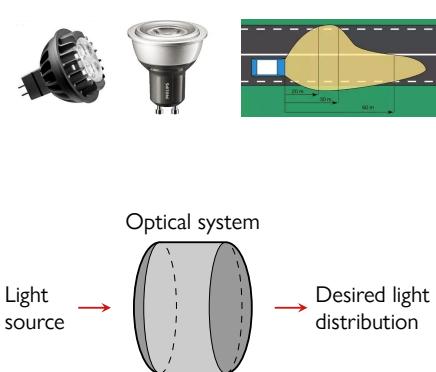
- ▶ LED lighting
- ▶ Road lights
- ▶ Car headlights

Industry standard: ray tracing

- ▶ Easy to implement
- ▶ Slow convergence
- ▶ Design, ray trace, change design, ray trace, ...

Inverse methods

- ▶ Directly compute required optical system
- ▶ Need solving PDE of Monge-Ampère type
- ▶ Avoid iterations and manual optimization

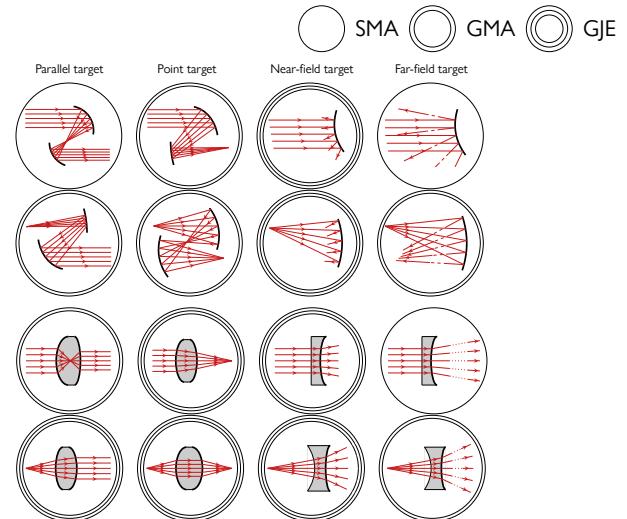


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Sixteen basic systems

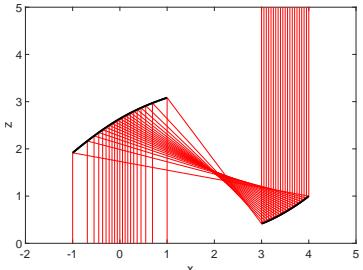
Reflector systems



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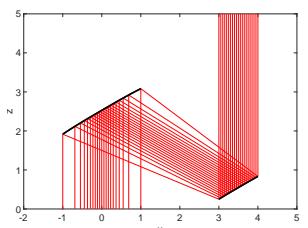
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Parallel-to-parallel reflector 2D



- Source: parallel beam with Gaussian light distribution
- Desired at target: parallel beam with uniform distribution
- Find the two freeform reflector surfaces

Energy conservation



- $S = (-1, 1)$, $T = (3, 4)$
- Source light distribution: $E = E(x)$, $x \in S$, Gaussian
- Target light distribution: $G = G(y)$, $y \in T$, uniform
- Energy conservation:
$$\int_{-1}^x E(\xi) d\xi = \pm \int_{m(-1)}^{m(x)} G(y) dy = \pm \int_{-1}^x G(m(\xi)) m'(\xi) d\xi$$
- Differentiate to x :
$$m'(x) = \pm \frac{E(x)}{G(m(x))}$$
- Solve ODE for m
- Differentiate $u(x) + w(y) = c(x, y)$ to x , substitute $y = m(x)$:
$$u'(x) = \frac{\partial c}{\partial x}(x, m(x))$$
- Solve ODE for u
- Second reflector: $w(m(x)) = c(x, m(x)) - u(x)$

Mathematical model: optimal transport formulation

Path of a ray

- Leaves source S at $P = (x, 0)$
- Hits first reflector at $A = (x, u(x))$
- Hits second reflector at $B = (y, L - w(y))$
- Arrives at target T at $Q = (y, L)$

Optical path length

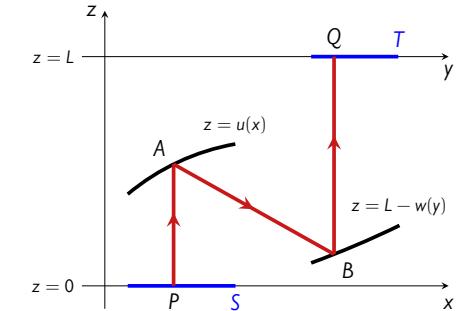
$$V = u(x) + d(A, B) + w(y)$$

- Let $\beta := V - L$. Then

$$\begin{aligned} u(x) + w(y) &= \frac{1}{2}\beta + L - \frac{(y-x)^2}{2\beta} \\ &=: c(x, y) \end{aligned}$$

$c(x, y)$ is called cost function

- Differentiate to x : **optical mapping** $y = m(x) = x + \beta u'(x)$



Parallel-to-parallel reflector, 2D and 3D

2D

Optical mapping:

$$y = m(x) = x + \beta u'(x) = \phi'(x)$$

Optimal transport formulation:

$$\begin{aligned} u(x) + w(y) &= c(x, y) \\ c(x, y) &= \frac{1}{2}\beta + L - \frac{(y-x)^2}{2\beta} \end{aligned}$$

Energy conservation:

$$m'(x) = \phi''(x) = \pm \frac{E(x)}{G(m(x))}$$

$$3D: \quad \mathbf{x} = \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} \quad \mathbf{y} = \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} \quad \mathbf{m} = \begin{pmatrix} m_1 \\ m_2 \end{pmatrix}$$

Optical mapping:

$$\mathbf{y} = \mathbf{m}(\mathbf{x}) = \mathbf{x} + \beta \nabla u(\mathbf{x}) = \nabla \phi(\mathbf{x})$$

Optimal transport formulation:

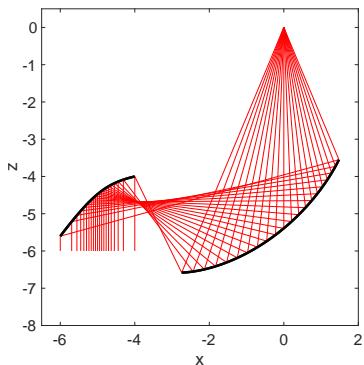
$$\begin{aligned} u(\mathbf{x}) + w(\mathbf{y}) &= c(\mathbf{x}, \mathbf{y}) \\ c(\mathbf{x}, \mathbf{y}) &= \frac{1}{2}\beta + L - \frac{\|\mathbf{y} - \mathbf{x}\|^2}{2\beta} \end{aligned}$$

Energy conservation:

$$\begin{aligned} \det(D\mathbf{m}(\mathbf{x})) &= \frac{\partial^2 \phi}{\partial x_1^2} \frac{\partial^2 \phi}{\partial x_2^2} - \left(\frac{\partial^2 \phi}{\partial x_1 \partial x_2} \right)^2 \\ &= \pm \frac{E(\mathbf{x})}{G(\mathbf{m}(\mathbf{x}))} \end{aligned}$$



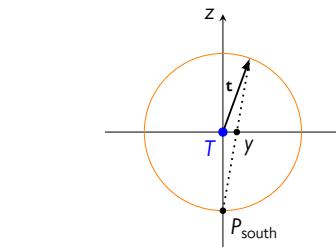
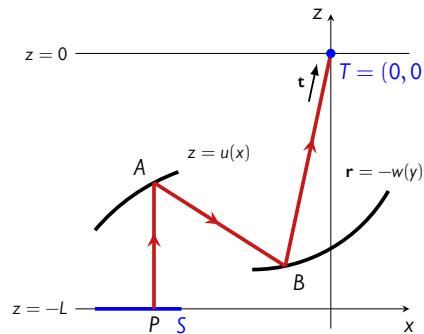
Parallel-to-point reflector 2D



- Source: parallel beam with Gaussian light distribution
- Point target with uniform distribution
- Find the two freeform reflector surfaces



Stereographic projection from the south pole



- Circle has radius 1, so $\mathbf{t} = \begin{pmatrix} t_1 \\ t_2 \end{pmatrix}$ $t_1^2 + t_2^2 = 1$
- South pole: $P_{\text{south}} = (0, -1)$
- One can show that

$$t_1 = \frac{2y}{1+y^2} \quad t_2 = \frac{1-y^2}{1+y^2} \quad y = \frac{t_1}{1+t_2}$$

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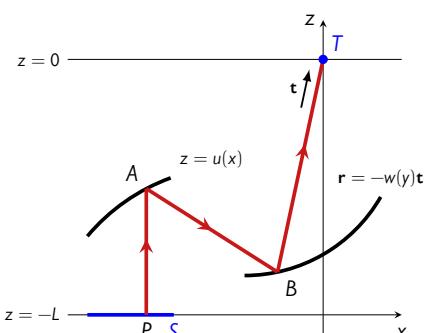
Mathematical model: optimal transport formulation



- Path of a ray
 - Leaves source S at $P = (x, -L)$
 - Hits first reflector at $A = (x, -L + u(x))$
 - Hits second reflector at $B = (-w(y)t_1, -w(y)t_2)$
 - Arrives at target $T = (0, 0)$
- Optical path length: $V = u(x) + d(A, B) + w(y)$
- Geometric relation

$$\left(-\frac{u}{\beta} - \frac{x^2}{2\beta^2} + \frac{V+L}{2\beta} \right) \cdot \left(\frac{\beta}{w} (1+y^2) - 2y^2 \right) = \left(1 + \frac{xy}{\beta} \right)^2$$

- Take logarithms: $u_1(x) + u_2(y) = c(x, y)$
Cost function is non-quadratic



Numerical method

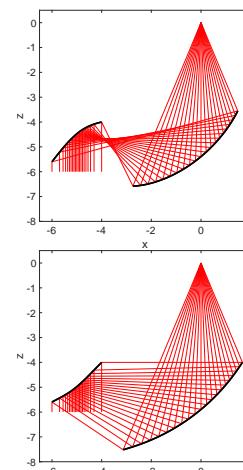


- Energy conservation:

$$m'(x) = \pm \frac{E(x)}{G(m(x))} \frac{1 + (m(x))^2}{2}$$
 Solve ODE for m
- Differentiate $u_1(x) + u_2(y) = c(x, y)$ to x , substitute $y = m(x)$:

$$u'_1(x) = \frac{\partial c}{\partial x}(x, m(x))$$
 Solve ODE for u_1
From u_1 compute u
- Second reflector:

$$u_2(m(x)) = c(x, m(x)) - u_1(x)$$
 From u_2 compute w



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Parallel-to-point reflector, 2D and 3D

2D

- Energy conservation:

$$m'(x) = \pm \frac{E(x)}{G(m(x))} \frac{1 + (m(x))^2}{2}$$

- Optimal transport formulation:

$$u_1(x) + u_2(y) = c(x, y)$$

$$c(x, y) = \log \left(\left(1 + \frac{xy}{\beta} \right)^2 \right)$$

3D: $\mathbf{x} = \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} \quad \mathbf{y} = \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} \quad \mathbf{m} = \begin{pmatrix} m_1 \\ m_2 \end{pmatrix}$

- Energy conservation:

$$\det(D\mathbf{m}(\mathbf{x})) = \pm \frac{E(\mathbf{x})}{G(\mathbf{m}(\mathbf{x}))} \frac{(1 + \|\mathbf{m}(\mathbf{x})\|^2)^2}{4}$$

- Optimal transport formulation:

$$u_1(\mathbf{x}) + u_2(\mathbf{y}) = c(\mathbf{x}, \mathbf{y})$$

$$c(\mathbf{x}, \mathbf{y}) = \log \left(\left(1 + \frac{\mathbf{x} \cdot \mathbf{y}}{\beta} \right)^2 \right)$$

- Differentiate, substitute $\mathbf{y} = \mathbf{m}(\mathbf{x})$, differentiate

$$\underbrace{D_{xy}c(\mathbf{x}, \mathbf{m}(\mathbf{x}))}_{\mathbf{C}} D\mathbf{m}(\mathbf{x}) = \underbrace{D^2u_1(\mathbf{x}) - D_{xx}c(\mathbf{x}, \mathbf{m}(\mathbf{x}))}_{\mathbf{P}}$$

- Generalized Monge-Ampère equation

$$\det(D\mathbf{m}(\mathbf{x})) = \frac{\det(\mathbf{P})}{\det(\mathbf{C})} = \pm \frac{E(\mathbf{x})}{G(\mathbf{m}(\mathbf{x}))} \frac{(1 + \|\mathbf{m}(\mathbf{x})\|^2)^2}{4}$$

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Mathematical model



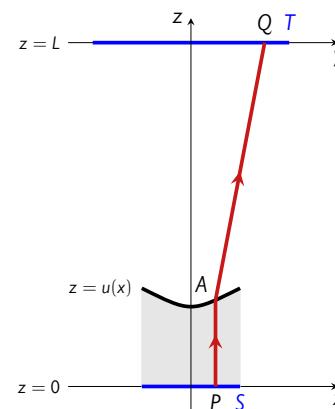
- Path of a ray

- Leaves source S at $P = (x, 0)$
- Hits freeform lens surface at $A = (x, u(x))$
- Arrives at target T at $Q = (y, L)$

- Optical path length:

$$\begin{aligned} V(y) &= n \cdot u(x) + d(A, Q) \\ &= n \cdot u(x) + \sqrt{(y-x)^2 + (L-u(x))^2} \end{aligned}$$

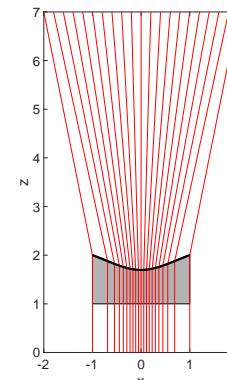
Cannot be formulated as optimal transport problem



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Parallel-to-near-field lens



- Source: parallel beam with Gaussian light distribution
- Near-field target with uniform distribution
- Find the single freeform lens surface

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Numerical method

Numerical method

- Energy conservation:

$$m'(x) = \pm \frac{E(x)}{G(m(x))}$$

Solve ODE for m

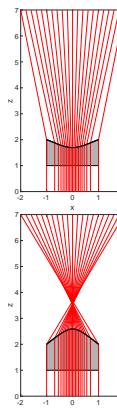
- Geometric relation:

$$V(y) = n \cdot u(x) + \sqrt{(y-x)^2 + (L-u(x))^2}$$

differentiate to x , substitute $y = m(x)$, solve for $u'(x)$:

$$u'(x) = \frac{m(x) - x}{n \cdot \sqrt{(m(x) - x)^2 + (L-u(x))^2} + u(x) - L}$$

Solve ODE for u



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Parallel-to-near-field lens, 2D and 3D

2D

- Energy conservation:

$$m'(x) = \pm \frac{E(x)}{G(m(x))}$$

- Geometric relation:

$$V(y) = n \cdot u(x) + \sqrt{(y-x)^2 + (L-u(x))^2}$$

3D

- Energy conservation:

$$\det(D\mathbf{m}(\mathbf{x})) = \pm \frac{E(\mathbf{x})}{G(\mathbf{m}(\mathbf{x}))}$$

- Geometric relation:

$$V(\mathbf{y}) = n \cdot u(\mathbf{x}) + \sqrt{\|\mathbf{y} - \mathbf{x}\|^2 + (L - u(\mathbf{x}))^2} \\ := H(\mathbf{x}, \mathbf{y}, u(\mathbf{x}))$$

- Differentiate to \mathbf{x} : $\nabla_{\mathbf{x}}(H(\mathbf{x}, \mathbf{y}, u(\mathbf{x}))) = \mathbf{0}$

- Let $\tilde{H}(\mathbf{x}, \mathbf{y}) := H(\mathbf{x}, \mathbf{y}, u(\mathbf{x}))$, substitute $\mathbf{y} = \mathbf{m}(\mathbf{x})$, differentiate to \mathbf{x} :

$$D_{xx}\tilde{H}(\mathbf{x}, \mathbf{m}(\mathbf{x})) + \underbrace{D_{xy}\tilde{H}(\mathbf{x}, \mathbf{m}(\mathbf{x}))}_{\mathbf{C}}(D\mathbf{m})(\mathbf{x}) = \mathbf{0}$$

- Generated Jacobian equation

$$\det(D_{xx}\tilde{H}(\mathbf{x}, \mathbf{m}(\mathbf{x}))) = \pm \det(\mathbf{C}) \frac{E(\mathbf{x})}{G(\mathbf{m}(\mathbf{x}))}$$



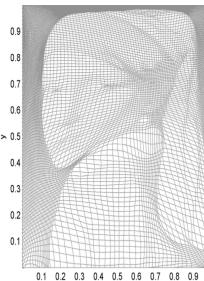
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Girl with a pearl earring (parallel-to-far-field reflector 3D)

- Parallel source
- Uniform source distribution
- Far-field target
- Single reflector



Target



Mapping



Ray trace

Prins, C. R., ten Thije Boonkamp, J. H. M., van Roosmalen, J., IJzerman, W. L., Tukker, T. W. (2014)
A Monge-Ampère-solver for freeform reflector design
SIAM Journal on Scientific Computing, 36(3), B640-B660
<https://doi.org/10.1137/130938876>



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Iterative least-squares solver for 3D systems



- Find mapping $\mathbf{m} : S \rightarrow T$ such that

$$\det(D\mathbf{m}) = \frac{E(\mathbf{x})}{G(\mathbf{m}(\mathbf{x}))}$$

$$\mathbf{m}(\partial S) = \partial T$$

- Break down in substeps
We compute \mathbf{P} , \mathbf{b} , \mathbf{m} such that

$$\mathbf{P} = \mathbf{C} D\mathbf{m}$$

$$\det(\mathbf{P}) = \det(\mathbf{C}) \frac{E(\mathbf{x})}{G(\mathbf{m}(\mathbf{x}))}$$

$$\mathbf{b}(\mathbf{x}) = \mathbf{m}(\mathbf{x}) \quad \mathbf{x} \in \partial S$$

\mathbf{b} maps ∂S to ∂T

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- Iterative procedure:

1. Choose an initial guess \mathbf{m}^0

Let $n = 0$

$$2. \text{ Let } J_l(\mathbf{m}, \mathbf{P}) = \frac{1}{2} \iint_S \|\mathbf{C} D\mathbf{m} - \mathbf{P}\|^2 d\mathbf{x}$$

$$\mathbf{P}^{n+1} = \operatorname{argmin}_{\mathbf{P}} J_l(\mathbf{m}^n, \mathbf{P})$$

Constrained minimization problem

$$3. \text{ Let } J_B(\mathbf{m}, \mathbf{b}) = \frac{1}{2} \int_{\partial S} \|\mathbf{m} - \mathbf{b}\|^2 ds$$

$$\mathbf{b}^{n+1} = \operatorname{argmin}_{\mathbf{b}} J_B(\mathbf{m}^n, \mathbf{b})$$

Projection on boundary ∂T

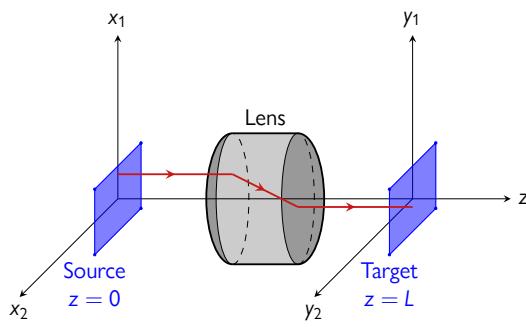
$$4. \text{ Let } J(\mathbf{m}, \mathbf{P}, \mathbf{b}) = \alpha J_l(\mathbf{m}, \mathbf{P}) + (1 - \alpha) J_B(\mathbf{m}, \mathbf{b})$$

$$\mathbf{m}^{n+1} = \operatorname{argmin}_{\mathbf{m}} J(\mathbf{m}, \mathbf{P}^{n+1}, \mathbf{b}^{n+1})$$

Elliptic PDEs for m_1 and m_2 — FVM

5. Let $n := n + 1$, go to Step 2

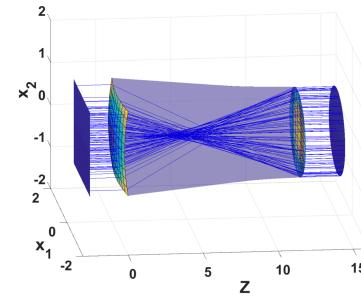
Laser beam shaping



- ▶ Light rays travel from left to right
- ▶ Source plane: $z = 0$
- ▶ Cartesian coordinates $\mathbf{x} = \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$
- ▶ Source S emits parallel light rays
- ▶ Exitance: $f(\mathbf{x}), \mathbf{x} \in S$
- ▶ Target plane: $z = L$
- ▶ Cartesian coordinates $\mathbf{y} = \begin{pmatrix} y_1 \\ y_2 \end{pmatrix}$
- ▶ Desired irradiance: $g(\mathbf{y}), \mathbf{y} \in T$
- ▶ First lens surface: $z = u(\mathbf{x}), \mathbf{x} \in S$
- ▶ Second lens surface: $L - z = w(\mathbf{y}), \mathbf{y} \in T$

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Numerical results for the laser beam shaping problem



Desired target distribution

Achieved

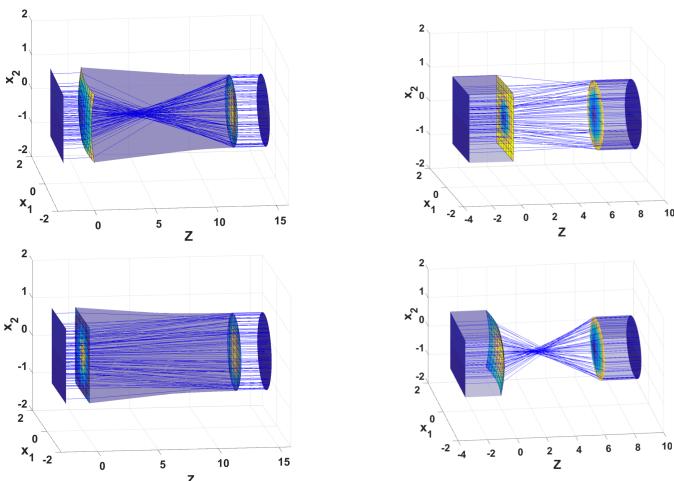


Yadav, N. K., ten Thije Boonkamp, J. H. M., IJzerman, W. L. (2019)
Computation of double freeform optical surfaces using a Monge–Ampère solver: Application to beam shaping
 Optics Communications, 439, 251-259
<https://doi.org/10.1016/j.optcom.2019.01.069>

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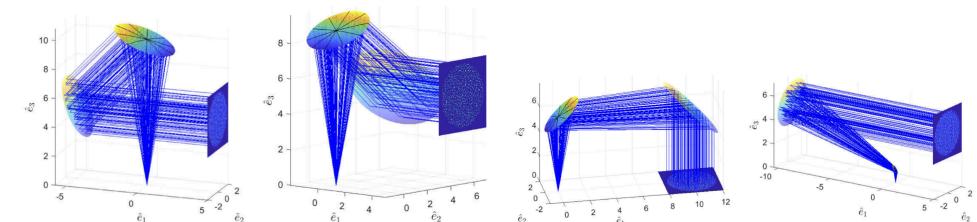
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Alternative optical systems



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Point-to-parallel reflector system 3D



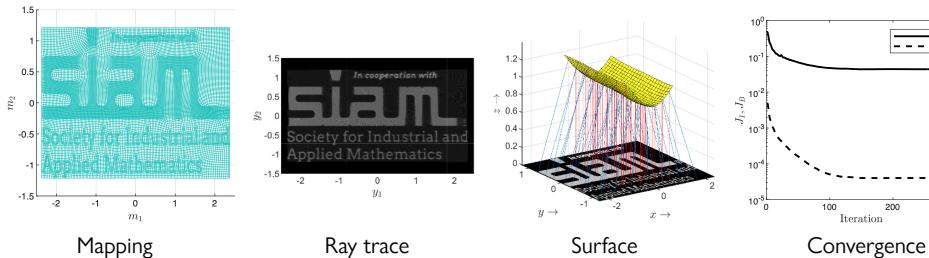
Van Roosmalen, A. H., Anthonissen, M. J. H., IJzerman, W. L., ten Thije Boonkamp, J. H. M. (2021)
Design of a freeform two-reflector system to collimate and shape a point source distribution
 Optics Express, 29(16), 25605-25625
<https://doi.org/10.1364/OE.425289>

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Parallel-to-near-field reflector 3D

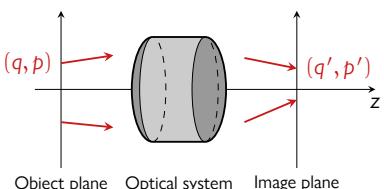


Romijn, L. B., Anthonissen, M. J. H., ten Thije Boonkamp, J. H. M., IJzerman, W. L. (2021)
An iterative least-squares method for generated Jacobian equations in freeform optical design
 SIAM Journal on Scientific Computing, 43(2), B298-B322
<https://doi.org/10.1137/20M1338940>

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Imaging optics

- Consider imaging optical system



- Each ray is defined by position q and direction p
- Hamiltonian optics:

$$\dot{q} = \frac{\partial H}{\partial p}, \quad \dot{p} = -\frac{\partial H}{\partial q}, \quad H = -\sqrt{n^2 - p^2}$$

- Optical map:

$$\begin{pmatrix} q' \\ p' \end{pmatrix} = M \begin{pmatrix} q \\ p \end{pmatrix}$$

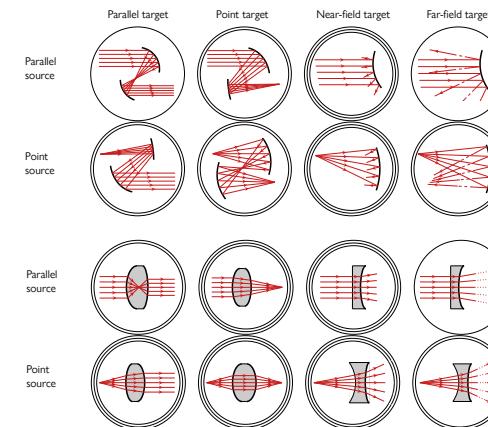
Ideal system:

$$M = M_{\text{linear}} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

- In practice M is nonlinear
- Goal: Mathematical description of nonlinearities, called **aberrations**

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Sixteen basic systems



- Model 1: Standard Monge-Ampère equation (SMA)**

Optimal transport formulation with quadratic cost function

- Model 2: Generalized Monge-Ampère equation (GMA)**

Optimal transport formulation with non-quadratic cost function

- Model 3: Generated Jacobian equation (GJE)**

No optimal transport formulation

Anthonissen, M. J. H., Romijn, L. B., ten Thije Boonkamp, J. H. M., IJzerman, W. L. (2021)
Unified mathematical framework for a class of fundamental freeform optical systems, Optics Express, 29(20), 31650-31664,
<https://doi.org/10.1364/OE.438920>

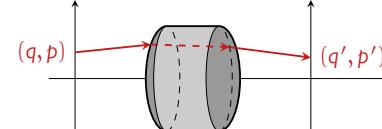
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Finding aberrations

- Consider imaging optical system



- Break down in steps: P—R—P—R—P

P: propagation

R: refraction

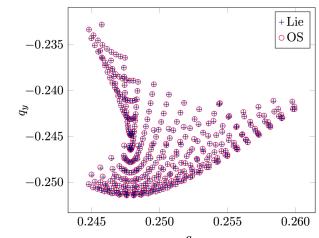
- Write each step as a **Lie transformation**

$$M_i = \exp([f_i, \cdot])$$

$$[f_i, \cdot] = \frac{\partial f_i}{\partial q} \frac{\partial}{\partial p} - \frac{\partial f_i}{\partial p} \frac{\partial}{\partial q}$$

- Optical map is product of M_i

- Spot diagram off-axis object

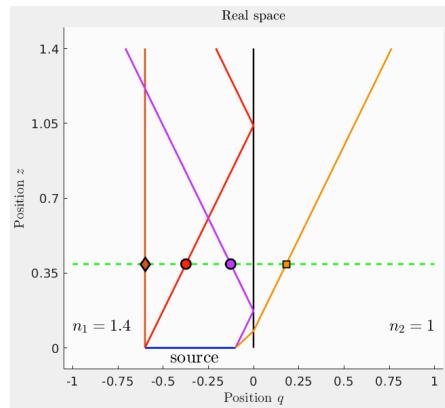


Barion, A., Anthonissen, M. J. H., ten Thije Boonkamp, J. H. M., IJzerman, W. L. (2023)
Computing aberration coefficients for plane-symmetric reflective systems: A Lie algebraic approach
 Journal of the Optical Society of America A, Optics, Image Science and Vision, 40(6), 1215-1224
<https://doi.org/10.1364/JOSAA.487343>

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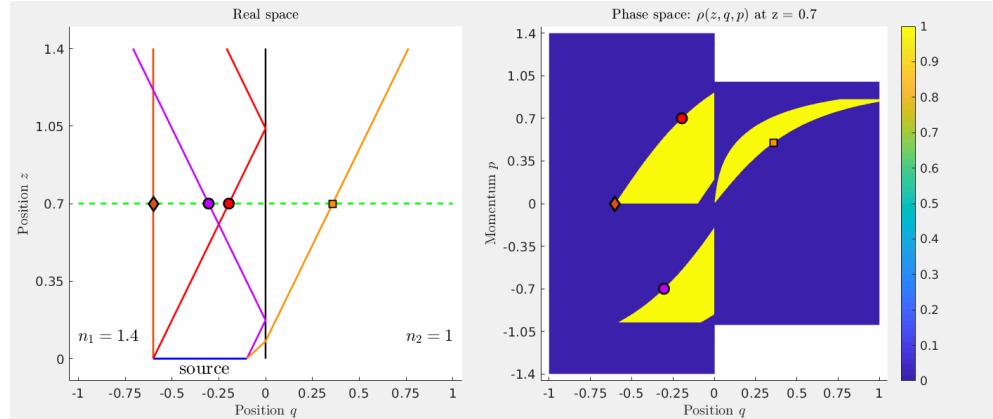
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Improved direct methods



- ▶ New methods work in **phase space**
- ▶ $q < 0$: Water, refractive index $n_1 = 1.4$
 $q > 0$: Air, $n_2 = 1$
 Light source at $z = 0$
- ▶ Each ray can be defined by position $q(z)$ and direction $p(z)$
- ▶ Four rays:
 - ◊ $p(0) = 0 \rightarrow q(z)$ constant
 - $p(0) > p_{\text{crit}}$ → ray is reflected
 - $p(0) > p_{\text{crit}}$ → ray is reflected
 - $0 < p(0) < p_{\text{crit}}$ → ray is refracted

Phase space



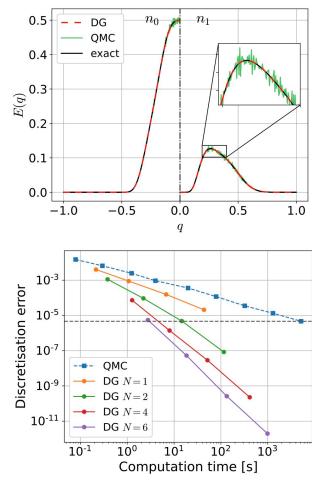
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Liouville's equation



- ▶ Conservation of luminance $\rho = \rho(z, q, p)$
- ▶ Liouville's equation
- ▶ Discontinuous Galerkin spectral element method (DG)
- ▶ Compared to quasi Monte Carlo ray tracing (QMC)

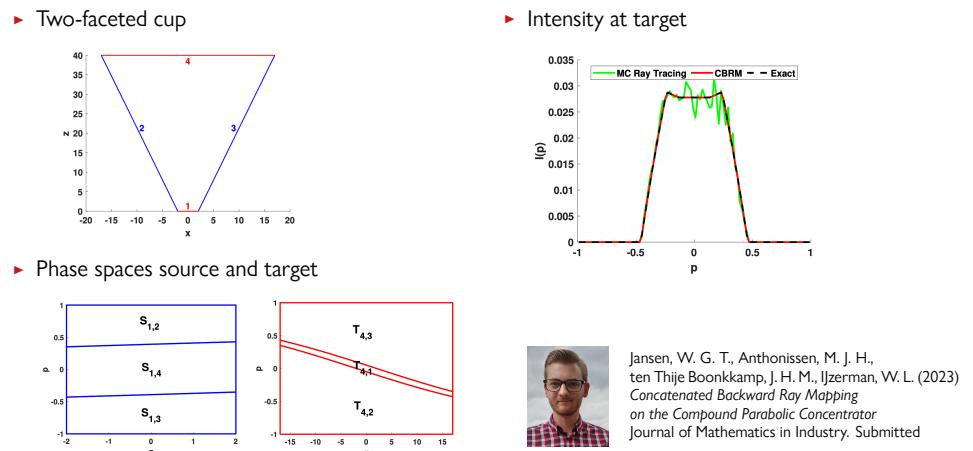


van Gestel, R. A. M., Anthonissen, M. J. H., ten Thije Boonkamp, J. H. M., IJzerman, W. L. (2023) An ADER discontinuous Galerkin method on moving meshes for Liouville's equation of geometrical optics. Journal of Computational Physics, 488, 112208 <https://doi.org/10.1016/j.jcp.2023.112208>

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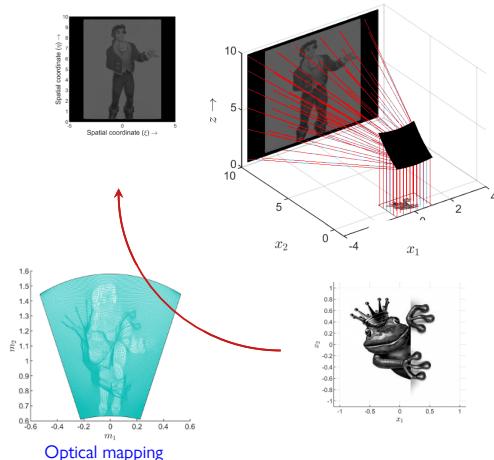
Backward ray mapping



Jansen, W. G. T., Anthonissen, M. J. H., ten Thije Boonkamp, J. H. M., IJzerman, W. L. (2023) Concatenated Backward Ray Mapping on the Compound Parabolic Concentrator. Journal of Mathematics in Industry. Submitted

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Conclusions

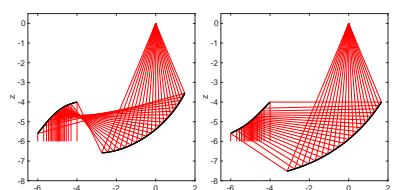


- Yes we can transform a frog into a prince with a mirror!
- And we can model and simulate another fifteen basic optical systems

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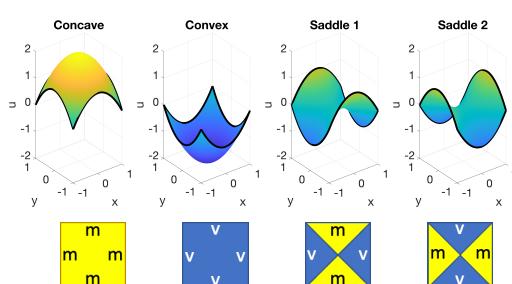
2D

$$m'(x) = \pm \frac{E(x)}{G(m(x))} \frac{1 + (m(x))^2}{2}$$



3D

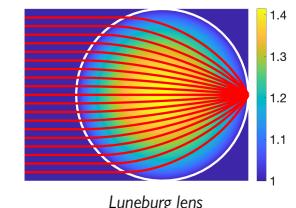
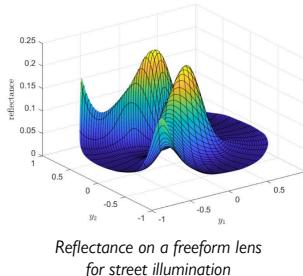
$$\det(D\mathbf{m}(\mathbf{x})) = \pm \frac{E(\mathbf{x})}{G(\mathbf{m}(\mathbf{x}))} \frac{(1 + \|\mathbf{m}(\mathbf{x})\|^2)^2}{4}$$



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Extensions and future work

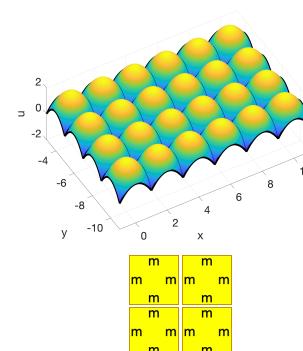
- Include Fresnel reflection and scattering
- Design lens arrays with smooth surfaces
- Use machine learning to speed up or replace the least-squares solver
- Model finite sources
- Model GRIN optics
- Move beyond the sixteen basic systems
- Include phase information



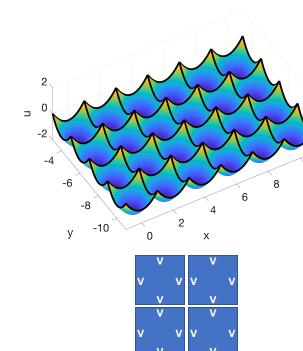
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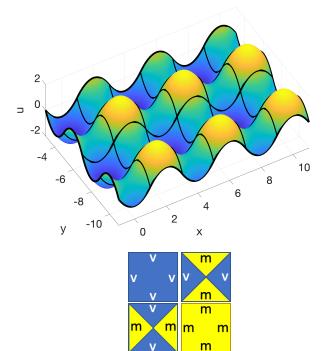
Lens arrays



Concave elements only
Not smooth at interfaces



Convex elements only
Not smooth at interfaces



Concave, convex and saddle
Smooth everywhere

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